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Two-crystal mid-infrared optical parametric oscillator for absorption and dispersion dual-comb spectroscopy

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We present a femtosecond optical parametric oscillator (OPO) containing two magnesium-doped periodically poled lithium niobate crystals in a singly resonant ring cavity, pumped by two mode-locked Yb-fiber lasers. As such, the OPO generates two idler combs (up to 220 mW), covering a wavelength range from 2.7 to 4.2 μm , from which a mid-infrared dual-comb Fourier transform spectrometer is constructed. By detecting the heterodyning signal between the two idler beams a full broadband spectrum of a molecular gas can be observed over 250 cm^{-1} within 70 μs with a spectral resolution of 15 GHz. The absorption and dispersion spectra of acetylene and methane have been measured around 3000 cm^{-1} , indicating that this OPO represents an ideal broadband mid-infrared source for fast chemical sensing. © 2014 Optical Society of America

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In the mid-infrared wavelength region between 2 and 20 μm , gases, liquids, and solids have unique absorption and dispersion features that are associated with their molecular rotational-vibrational transitions. To observe these spectra, a wide variety of instrumentations is available for quantitative analysis, such as Fourier transform spectrometry (FTS). With the recent development of laser frequency combs, a novel approach can be used for measuring the broadband spectra through multiheterodyne spectroscopy [1].

To date, commercially available frequency combs are limited to the near-infrared region up to 2 μm . Therefore, spectrometers based on frequency combs offer limited sensitivity, as the molecular transitions have typically two orders of magnitude weaker absorption cross sections in the near-infrared wavelength region as compared to the mid-infrared. This is why research on mid-infrared frequency comb sources is of great interest. Four strategies can be considered: the use of mode-locked lasers with a gain medium producing direct mid-infrared radiation, different frequency generation (DFG), optical parametric oscillation, or Kerr comb generation [2]. Among those, the optical parametric oscillator (OPO) offers a wide tuning range in the mid-infrared region and high conversion efficiency. In the past, OPOs pumped by Ti:sapphire femtosecond lasers were demonstrated as efficient sources for broadband and sensitive spectroscopy [3,4]. Later on, fully stabilized OPO-based combs have been demonstrated for wavelengths longer than 2.5 μm , either via singly resonant OPOs [5,6] or double-resonant OPOs [7,8] generating ultrabroad mid-infrared spectra.

Different approaches have been developed to use frequency combs systems as light sources for absorption spectroscopy. Taking advantage of the wide spectral emission, near-infrared molecular spectra were recorded with multichannel arrangements [9–11], multiplex spectrometers [12,13] or different types of dual-comb spectrometry [1,14–19]. In combination with a multipass gas cell or a high finesse optical cavity, it is possible to

achieve high sensitivity to trace molecular gases at part per billion levels [1,9,10,20,21].

The development of dual-comb FTS overcomes several disadvantages of conventional FTS [14,22] and gives simultaneous measurements of the absorption and dispersion spectra of a sample. Several approaches have been demonstrated to perform mid-infrared dual-comb spectroscopy [16,18,23]. Among those, DFG-based systems have provided accurate spectral measurements around 3.4 μm [16]. However, microwatt power levels and the relative narrow spectral bandwidth of the DFG source can limit its implementation in applications such as real-time multigas detection.

Here, we report on a high-power mid-infrared OPO containing two MgO-doped periodically poled lithium niobate (MgO:PPLN) crystals in a singly resonant ring cavity, pumped by two mode-locked lasers. If the idea of two-crystal OPOs has been demonstrated for the generation of continuous wave near-infrared light [24], for the first time it is implemented on mid-infrared frequency combs. The special and unique design of the OPO cavity offers several remarkable advantages for dual-comb spectroscopy, such as the emission of two independent mid-infrared beams separated in space and a significant average optical power of 220 mW.

The configuration of the OPO is depicted in Fig. 1. The pump lasers (Menlo Systems, Germany) are femtosecond Yb-fiber lasers, emitting at a central wavelength of 1040 nm (9615 cm^{-1}), delivering up to 2 W of average power with 80 fs optical pulse duration (30 nm or 275 cm^{-1} spectral width at 1/e value); the collimated beams have diameters of 2.4 and 2.8 mm. The repetition frequency of one pump laser (90 MHz) is stabilized by synchronizing its 10th harmonic to a frequency generator (R&S, SMB100A) offering a single sideband phase noise of -128 dBc at 1 GHz. The 10th harmonic of the repetition rate of the second pump laser is locked to the same reference clock with a slightly different frequency of up to 3 kHz, corresponding to a 300 Hz different repetition

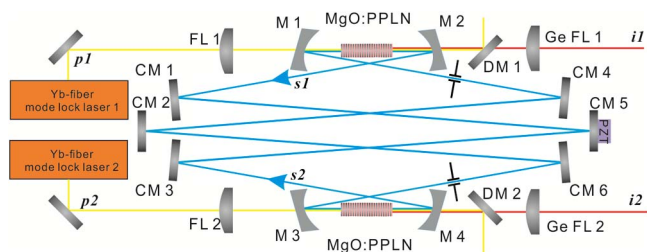


Fig. 1. Experimental setup of the OPO cavity. p1–2, s1–2, i1–2, pumps, signals, and idler beams, respectively, corresponding to different pump lasers. FL1–2, focal lenses ($f = 100$ mm). M1–M4, curved mirrors ($r = 100$ mm). CM1–CM6, chirped mirrors (100 fs² at 1500 nm). DM1–2, dichroic mirrors. GeFL1–2, germanium focal lenses ($f = 100$ mm). PZT, piezo actuator.

frequency. A higher difference in between the repetition frequencies will not give simultaneous oscillation within the OPO cavity for both beams. The pump beams are focused onto the OPO crystals inside the ring cavity with two lenses ($f = 100$ mm). The 5 mm long MgO:PPLN crystals (Covesion, UK) are temperature stabilized and have eight poling periods each, ranging from 27.91 to 31.59 μm . Both end facets of each PPLN crystal are anti-reflective (AR) coated ($R < 1.5\%$ at 1064 nm, $R < 1\%$ at 1400 – 1800 nm, and $R < 6\%$ at 2600 – 4800 nm). The OPO cavity is designed to be synchronously pumped; therefore, the cavity length is about 3.4 m long. The ring cavity consists of four curved mirrors (radius of curvature 100 mm, AR coated at 1064 and 3650 – 4850 nm and high-reflective [HR] coated at 1350 – 1500 nm) and six flat, chirped mirrors (HR coated at 1375 – 1750 nm, Layertec GmbH, Germany). The chirped mirrors (CM1–CM6) have a second-order dispersion of approximately -100 fs² per bounce at 1500 nm. They compensate the group velocity dispersion introduced by the PPLN crystals and the curved mirrors at 1.5 μm . A piezo actuator is mounted on CM5 for further frequency combs stabilization. The idler beams are emitting collinearly with the residual pump beams, and two dichroic mirrors (LaserOptik GmbH, $R > 99.9\%$ at 1064 nm and $R < 1\%$ at 3200 – 3900 nm) are used to separate the pump and idler beams. Since a wide variety of non-phase-matched light beams are generated from the PPLN crystals, two germanium lenses ($f = 100$ mm, AR coated at 3 – 5 μm , Edmund Optics Ltd, UK) are filtering out and collimating the idler beams. To impede higher order transversal modes in the OPO cavity, two irises are placed. The aperture of the irises causes significant diffraction losses for the higher-order modes. This is beneficial for heterodyne detection since the two output idler beams should have an optimal spatial overlap, which can be achieved when both beams are in TEM₀₀ mode [25].

The spectral properties of the OPO are analyzed by a rapid-scan Fourier transform spectrometer (Magna-IR E.S.P system). For this experiment, the OPO cavity is pumped by only one mode-locked laser. The spectra of the signal beam for the eight poling periods of the crystal are depicted in Fig. 2(a) (spectral resolution of 0.4 cm⁻¹), ranging from 4800 to 7200 cm⁻¹ (1.3 – 2.2 μm). The idler spectra are depicted in Fig. 2(b). As can be seen from the figure, a spectral range extending from 2300 to 3600 cm⁻¹ (2.7 – 4.4 μm) can be covered. The maximum spectral

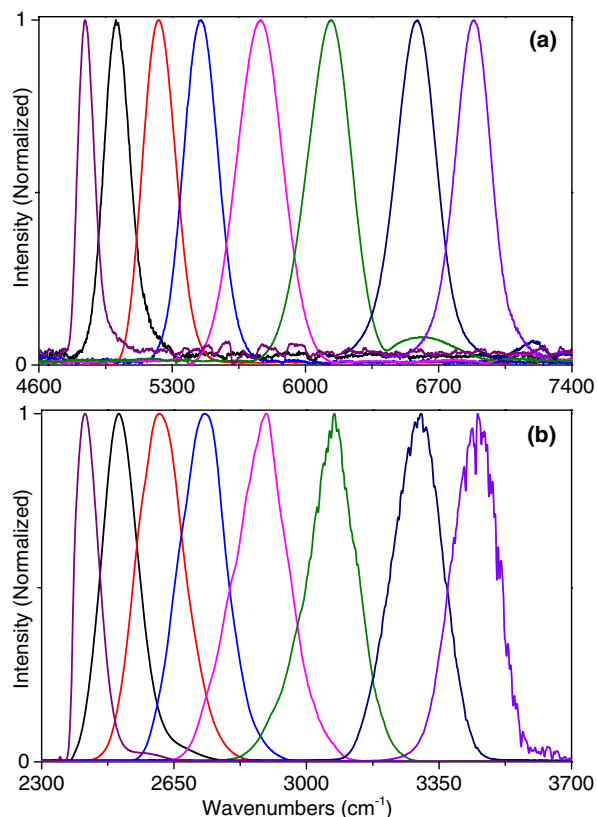


Fig. 2. (a) Spectral intensity distribution of the signal using eight poling periods, 27.91 to 31.59 μm , covering a wavelength region between 4800 and 7200 cm⁻¹. (b) Spectral intensity distribution of the idler light using eight poling periods, resulting in a wavelength coverage between 2400 and 3600 cm⁻¹.

width within one period is 300 cm⁻¹ ($1/e$ value), centered at 3080 cm⁻¹ (3.24 μm). To characterize the threshold and conversion efficiency, the idler power has been measured at different crystal periods for different pump powers (see Fig. 3). A pump power threshold as low as 46 mW has been achieved for the poling period generating an idler spectra centered at 3300 cm⁻¹ with a maximum conversion efficiency of 27% .

To conduct spectral analysis of gas phase samples, a dual-comb FTS is constructed, which is shown in Fig. 4.

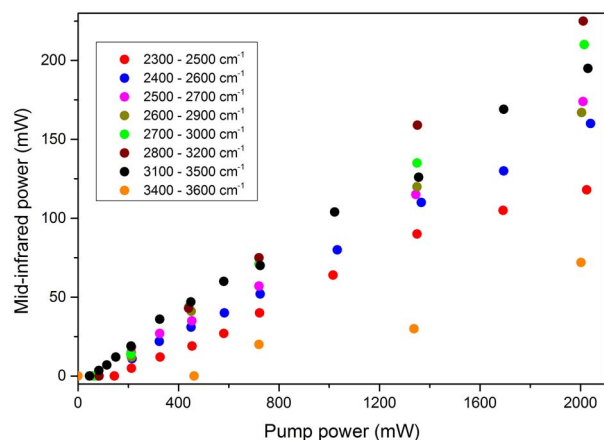


Fig. 3. Average output power of the idler at the eight different crystal periods.

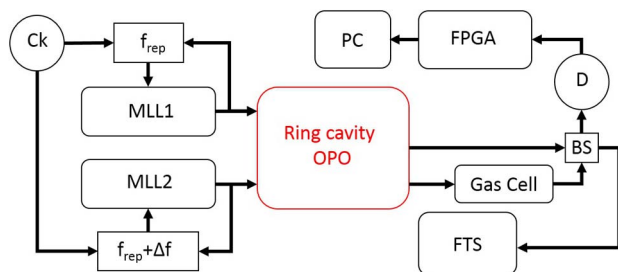


Fig. 4. Experimental setup of the dual-comb FTS. Two Yb-fiber mode-locked lasers (MLL1–2) are coupled to a singly resonant ring cavity OPO. Synchronized to a reference clock (Ck), the difference between the repetition frequencies is Δf . The OPO generates two spatially separated idler beams. One idler beam probes the gas sample in a 30 cm long single pass cell (Gas Cell) and is recombined with the second idler beam at the beam splitter (BS). Both overlapped beams are focused and the beating signal is detected by a Peltier cooled fast IR detector (D). The FFT of the detector signal is calculated by a FPGA, giving a spectral image of the gas sample. The light from the other side of the beam splitter is analyzed by a commercial rapid-scan FTS.

It consists of two femtosecond Yb-fiber lasers, each pumping the singly resonant OPO cavity described previously. The repetition rates of both pump lasers are synchronized to the same frequency generator. The difference in repetition rate (Δf) between the two lasers is set to values ranging from 40 to 120 Hz for various experiments. One collimated idler beam is probing the gas sample before interfering with the other idler beam on a single detector (Vigo, PVI-4TE, 50 MHz bandwidth). From this fast preamplified Peltier cooled mid-infrared photo-detector, the interferogram is low-pass filtered (up to 50 MHz) and sampled at a frequency of 125 MHz. A fast Fourier transform (FFT) of the recorded signal, calculated with a field programmable gate array (FPGA)–based data acquisition system (National Instruments Nederland BV) gives a real-time downconverted image of the molecular spectra in the RF frequency domain.

Figure 5 gives an illustration of the truncated interferogram from the full recording and the calculated spectra

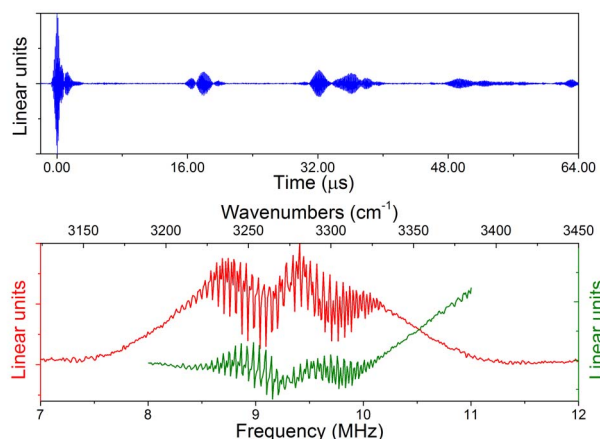


Fig. 5. Time domain interferogram (upper panel) of acetylene (C_2H_2 , 10% in N_2 , pressure 100 mbar, 30 cm path length, $\Delta f = 46$ Hz). Lower panel: Power spectrum (red) and dispersion information (green) in the RF domain obtained by considering the 120 μs interferogram.

of acetylene (C_2H_2) at a concentration of 10% in nitrogen (pressure 100 mbar, 30 cm path length, $\Delta f = 46$ Hz). On the interferogram, the wave packages after the initial burst indicate that optical frequency components of the comb source are absorbed and dispersed by the gas sample. The FFT of the interferogram recorded within 120 μs gives the acetylene absorption spectrum between 3150 and 3400 cm^{-1} with a RF resolution of 7.6 kHz, corresponding to an infrared spectral resolution of 15 GHz (0.5 cm^{-1}). Figure 6(a) shows the first 12 μs of the recorded interferogram of methane (CH_4 , 1% in nitrogen, 1 bar pressure, 30 cm path length, $\Delta f = 117$ Hz). The recorded 60 μs interferogram carries the information of a 250 cm^{-1} wide spectrum between 2950 and 3200 cm^{-1} with a RF resolution of 15.2 kHz, corresponding to an infrared spectral resolution of 11.7 GHz (0.39 cm^{-1}). In Fig. 6(c), the absorption and dispersion spectra from methane are obtained from a single shot interferogram within 60 μs and compared after background removal [26,27] with a simulated spectrum, based on the HITRAN database (0.4 cm^{-1} resolution). The experimental and simulated spectra agree quite well for the recorded P, Q and R branches. Below 3000 cm^{-1} , distortions are present due to a lower signal-to-noise ratio (SNR) observed in the RF spectrum (below 30 MHz) enhanced by the background removal calculation [26]. The other spectral distortions are mainly due to the remaining frequency jittering of the laser combs since the offset frequencies of the pumps are free-running [28].

The OPO cavity reported has several advantages for dual-comb spectroscopy. Both generated idlers originate from the same OPO cavity experiencing identical disturbances, such as mirror vibrations and air or temperature fluctuations. The shared OPO cavity cancels out much of the instabilities compared to using two separate OPOs as dual-comb sources. Moreover, the two signal beams resonating in the optical cavity are counterpropagating and

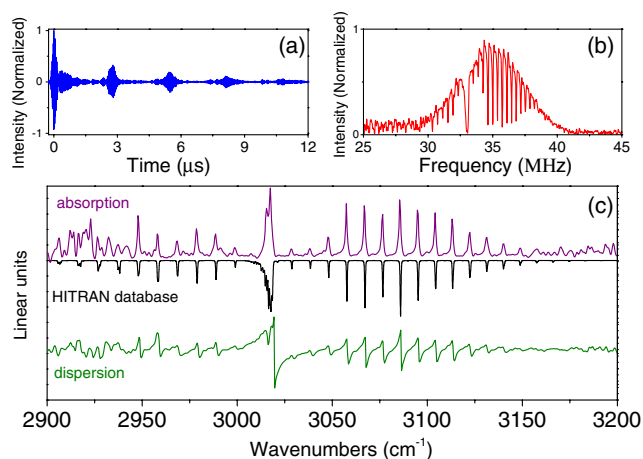


Fig. 6. (a) Time domain interferogram of methane (CH_4) over 12 μs . (b) Apodized power spectrum calculated from a 60 μs recorded interferogram. (c) Comparison between the measured (purple curve) and simulated (black curve, HITRAN database) absorption spectrum of methane. The Q branch (at 3050 cm^{-1}), the R branch (right side), and the beginning of the P branch (left side) of the ν_3 vibrational transition of methane can be clearly observed. The green curve shows the dispersion information after background removal.

do not affect the generation of the two idler beams. With this new cavity design, the spectrum overlap of the two idlers can be maximized by independently adjusting the temperature of each crystal. Furthermore, by two independent outputs, dispersion spectra can be measured, which might be beneficial for trace gas sensing [29,30]. For the dual-comb configuration as well as for Fourier spectroscopy, the amplitude of the dispersion signal should be immune to the intensity fluctuations and can provide more reliable and accurate background free data than with the absorption spectrum [26,27]. In our system, the main remaining noise sources are the detector noise and the frequency jittering of the combs (due to the pump fluctuations). They affect the SNR, the spectral resolution, and the efficiency of long-time averaging. A better stabilization of the combs or adaptive sampling techniques [28] will be considered to achieve measurement with higher spectral resolution and higher SNR by multiple spectrum averages.

In summary, we have reported a dual-comb FTS based on a high-power singly resonant OPO with two MgO:PPLN crystals in a ring cavity. This is demonstrated with an absorption and dispersion spectrum of methane covering 250 cm^{-1} ($2950\text{--}3200\text{ cm}^{-1}$) in the mid-infrared wavelength region, recorded within $60\text{ }\mu\text{s}$. In addition, the acetylene spectrum is recorded by tuning the OPO to another spectral region covering 250 cm^{-1} and centered at 3250 cm^{-1} . Since the mid-infrared frequency combs are not yet fully stabilized, these initial results give an overview of the advantages offered by the presented innovative OPO cavity to achieve accurate, high-spectral-resolution, broadband absorption and dispersion spectroscopy for real-time chemical sensing.

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References

1. I. Coddington, W. C. Swann, and N. R. Newbury, *Phys. Rev. Lett.* **100**, 013902 (2008).
2. A. Schliesser, N. Picque, and T. W. Hänsch, *Nat. Photonics* **6**, 440 (2012).
3. K. A. Tillman, R. R. J. Maier, D. T. Reid, and E. D. McNaghten, *Appl. Phys. Lett.* **85**, 3366 (2004).
4. J. H. Sun, B. J. S. Gale, and D. T. Reid, *Opt. Lett.* **32**, 1414 (2007).
5. F. Adler, K. C. Cossel, M. J. Thorpe, I. Hartl, M. E. Fermann, and J. Ye, *Opt. Lett.* **34**, 1330 (2009).
6. Z. Zhang, X. Fang, T. Gardiner, and D. T. Reid, *Opt. Lett.* **38**, 2077 (2013).
7. N. Leindecker, A. Marandi, R. L. Byer, and K. L. Vodopyanov, *Opt. Express* **19**, 6296 (2011).
8. N. Leindecker, A. Marandi, R. L. Byer, K. L. Vodopyanov, J. Jiang, I. Hartl, M. Fermann, and P. G. Schunemann, *Opt. Express* **20**, 7046 (2012).
9. M. J. Thorpe and J. Ye, *Appl. Phys. B* **91**, 397 (2008).
10. S. A. Diddams, L. Hollberg, and V. Mbele, *Nature* **445**, 627 (2007).
11. C. Gohle, B. Stein, A. Schliesser, T. Udem, and T. W. Hänsch, *Phys. Rev. Lett.* **99**, 263902 (2007).
12. F. Adler, P. Masłowski, A. Foltynowicz, K. C. Cossel, T. C. Briles, I. Hartl, and J. Ye, *Opt. Express* **18**, 21861 (2010).
13. J. Mandon, G. Guelachvili, and N. Picque, *Nat. Photonics* **3**, 99 (2009).
14. F. Keilmann, C. Gohle, and R. Holzwarth, *Opt. Lett.* **29**, 1542 (2004).
15. A. Schliesser, M. Brehm, F. Keilmann, and D. van der Weide, *Opt. Express* **13**, 9029 (2005).
16. E. Baumann, F. R. Giorgetta, W. C. Swann, A. M. Zolot, I. Coddington, and N. R. Newbury, *Phys. Rev. A* **84**, 062513 (2011).
17. T. Ideguchi, A. Poisson, G. Guelachvili, T. W. Hänsch, and N. Picqué, *Opt. Lett.* **37**, 4847 (2012).
18. Z. Zhang, T. Gardiner, and D. T. Reid, *Opt. Lett.* **38**, 3148 (2013).
19. B. Bernhardt, A. Ozawa, P. Jacquet, M. Jacquety, Y. Kobayashi, T. Udem, R. Holzwarth, G. Guelachvili, T. W. Hänsch, and N. Picque, *Nat. Photonics* **4**, 55 (2009).
20. M. J. Thorpe, K. D. Moll, R. J. Jones, B. Safdi, and J. Ye, *Science* **311**, 1595 (2006).
21. M. J. Thorpe, D. Balslev-Clausen, M. S. Kirchner, and J. Ye, *Opt. Express* **16**, 2387 (2008).
22. S. Schiller, *Opt. Lett.* **27**, 766 (2002).
23. Y. Wang, M. G. Soskind, W. Wang, and G. Wysocki, *Appl. Phys. Lett.* **104**, 031114 (2014).
24. G. K. Samanta, A. Aadhi, and M. Ebrahim-Zadeh, *Opt. Express* **21**, 9520 (2013).
25. P. Vladimir, *Laser Heterodyning* (Springer, 2009).
26. J. A. de Haseth and P. R. Griffiths, *Fourier Transform Infrared Spectrometry* (Wiley, 2007).
27. J. Kauppinen and J. Partanen, *Fourier Transforms in Spectroscopy* (Wiley, 2001).
28. T. Ideguchi, A. Poisson, G. Guelachvili, N. Picqué, and T. W. Hänsch, *Nat. Commun.* **5**, 3375 (2014).
29. G. Wysocki and D. Weidmann, *Opt. Express* **18**, 26123 (2010).
30. M. Nikodem and G. Wysocki, *Ann. N.Y. Acad. Sci.* **1260**, 101 (2012).